



On the origin of the Cold Spot Kaiki Taro Inoue (KINDAI)

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Introduction

The CMB Cold Spot ~3σ deviation in kurtosis using spherical mexican-hat wavelets (Cruz et al. 2003)





The CMB Cold Spot

- ~3σ deviation in kurtosis using spherical mexican-hat wavelets (Cruz et al. 2003)
- but not significant for non-compensated filters (Smith & Huterer 2010)
- ~3σ deviation for signals smoothed by a top-hat compensating filter (Inoue, Sakai & Tomita 2010)



Temperature profile of the Cold Spot

1.0



Fig. 10.— Left: the WMAP7 ILC temperature map $(40^{\circ} \times 40^{\circ})$ smoothed at 1° scale. Right: the averaged radial profile of the ILC map as a function of inclination angle θ from the center of the cold spot $(l, b) = (207.8^{\circ}, -56.3^{\circ})$. A peak at $\theta \sim 15^{\circ}$ corresponds to a hot ring.

Compensated top-hat filter



 θ_{in}

 $\theta_{_{out}}$

 θ $W(\theta) = \begin{cases} 1 & , \ \theta < \theta_{in} \\ -1 & , \ \theta_{in} \le \theta < \theta_{out} \end{cases}$ $\int_{0}^{\theta_{out}} 2\pi\theta W(\theta) \sin\theta \, d\theta = 0$

For those who like math ...

where $A = 2\pi (1 - \cos \theta_{in})$. Plugging ΔT expanded in spherical harmonics Y_{lm} ,

$$\Delta T = \sum_{lm} a_{lm} Y_{lm} \qquad (C2)$$

into equation (A1), we have

$$\Delta T_f = A^{-1} \sum_l a_{l0} G_l, \tag{C3}$$

where

$$G_{l} = \frac{\sqrt{\pi(2l+1)}}{l+1} \left[2\left(-x_{in}P_{l}(x_{in}) + P_{l-1}(x_{in})\right) + x_{out}P_{l}(x_{out}) - P_{l-1}(x_{out}) \right],$$

$$x_{in} = \cos\theta_{in}, \ x_{out} = 2x_{in} - 1.$$
(C4)

Note that we have used a formula for the Legendre function P_{l} ,

$$\frac{dP_l(x)}{dx} = \frac{l(l+1)}{1-x^2} \int_x^1 P_l(x) dx$$
(C5)

in deriving equation A4. Because ΔT is assumed to be isotropic on S^2 , the variance of ΔT_f can be written as a function of the angular power spectrum C_l as

$$\sigma_f^2 = A^{-2} \sum_l C_l G_l^2.$$
 (C6)

If the CMB sky is smoothed by a Gaussian beam with the FWHM θ_s , then the variance is $\sigma_f^2 = A^{-2} \sum_l C_l B_l G_l^2,$ (C7)

where $B_{\rm I} = \exp[-\sigma_s^2 l(l+1)/2]$, and $\sigma_s = (8 \ln 2)^{-1/2} \theta_s$.

Temperature profile of CS



Statistical significance



Statistical significance for the full sky

Sample number~observed area/size of filtering

At =5 degree, significance is 0.7-1 %(~2σ)
 at =12 degree, significance is 0.01-0.2%(~3σ)

Optimal filter for a cold spot surrounded by a hot ring?

Origin of the CS

- A supervoid with r~200Mpc/h, δ~-0.3 at z~1 (Inoue & Silk 2006)
- A hot ring cannot be produced by a compensating void (Tomita & Inoue 2007 Inoue & Sakai 2008)
- A texture at z~5 (Cruz et al. 2007)
- Statistical fluke?

Supervoid?

- A dip of 20-45% in the surface brightness in radio bands (Rudnick, Brown, & Williams 2007)
- No significant dip in galaxy counts at 0.3<z<1 in the l.o.s to the center of the CS (Bremer et al. 2010, Granett et al. 2010)
 - A low-density region at z~0.2 in 2MASS galaxy catalog may contribute 8µK to the CS (Francis & Peacock 2010)

Supervoids?



(Inoue & Silk, 2006&2007)

"the Cold Spot"



(Inoue & Silk 2007)

2MASS

Local ISW



Figure 1. (Left) The 2D reconstruction of the local density field described in Section 2.1 in three photometric redshift shells: 0.0 < z < 0.1 (top), 0.1 < z < 0.2 (middle) and 0.2 < z < 0.3 (bottom). The plots show overdensity δ on a scale $-0.6 \leq \delta \leq 0.6$. (Right) The corresponding ISW signal in mK computed from the reconstructed density field using equation (4).

0.0<Z<0.1

0.1<Z<0.2

0.2<Z<0.3 (Francis & Peacock, 2010)

Questions

Q1. If a void is non-compensated, can a hot ring be produced and statistically more probable?

A1. Yes it can. But not probable.

Q2. Even if it is a fluke, what is the expected ISW contribution from a possible low-density region to the CS?

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A2. It is about 10 %.

For those who want to know the reason...



ISW effect due to spherical voids/clusters

ISW (RS) effects from SSS

Theory

- Thin-shell approx. (Inoue & Silk 2006&2007)
 2nd order perturbations (Tomita & Inoue 2008)
 LTB solutions (Sakai & Inoue 2008)
 Observation
 - SDSS LRGs (Granett et al. 2008)
 2MASS photo-z (Francis & Peacock 2009)

Integrated Sachs-Wolfe (ISW) Effect temperature fluctuation due to time-evolving gravitational potential



→low temperature for CMB photons that pass through a void

temperature fluctuation due to time-evolving gravitational potential



→high temperature for CMB photons that pass through a cluster

Second order ISW (RS) Effect



Second order ISW (RS) Effect







1st & 2nd order ISW effects

cluster

void



(Tomita & Inoue 2008)

Non-compensation effect

Massive wall around a void could increase the CMB temperature.

Less massive "void wall" around a cluster could decrease the CMB temperature.

Quasi-linear perturbation can be estimated
 using 2nd order perturbation (Tomita & Inoue 2008)



Compensating $mod \varepsilon = 0$







Non-compensation effect void density contrast cluster



Non-compensation effect void temperature variation cluster



Non-compensation effect

- A hot ring surrounding a cold spot can be produced by an over-compensated void even in an accelerating epoch.
- A cold ring surrounding a hot spot can be produced by an under-compensated cluster even in an accelerating epoch.
- A cold ring surrounding a cold spot can be produced by an under-compensated non-linear cluster.

Non-compensation effect

- A hot ring surrounding a cold spot can be produced by an over-compensated void even in an accelerating epoch.
- A cold ring surrounding a hot spot can be produced by an under-uncompensated cluster even in an accelerating epoch.
- A cold ring surrounding a cold spot can be produced by an under-uncompensated non-linear cluster.

Late ISW effect & other CMB effects

Temperature anisotropy Ordinary Sachs-Wolfe (OSW) effect : spatial fluctuation on the last scattering surface. Doppler effect \triangleright :spatial fluctuation of the fluid velocity on LSS. Early integrated Sachs-Wolfe (ISW) effect \succ : time evolving potential at radiation->matter epoch. Late integrated Sachs-Wolfe (ISW) effect : time evolving potential at matter-> Λ epoch.



angular power ofsqrt \times 2.73K

Temperature anisotropy smoothed by a compensated top-hat filter



Statistical analysis

A single void model

We assume there is a spherical top-hat type void (low density region) in the l.o.s. to the CS at z<1.5.</p>

We assume a concordant cosmology with

$$(\Omega_{\Lambda} = 0.74, \Omega_m = 0.26, \Omega_b = 0.044, \sigma_8 = 0.80, h = 0.72, n = 0.96)$$

We assume that the CS is produced by a late ISW due to a void at z<1.5 (low density region) plus the OSW, earlyISW, and Doppler effects. No correlation with the fluctuations at the LSS is assumed.

Over-compensated void





 $\Delta T_{f}(1\sigma)$ $=18 \,\mu K$

ε=1

 $\theta_{\rm in}$ =12°

 $\theta_2 = 30^\circ$

Compensated void



 $\Delta T_{f}(1\sigma)$ =18 μ K $\epsilon=0$ $\theta_{in}=12^{\circ}$

$$\theta_1 = \frac{\theta_0 + \theta_2}{2}$$

$$\theta_2 = 30^\circ$$

Under-compensated void



 $\Delta T_{f}(1\sigma)$ =18 μ K ϵ =-1 θ_{in} =12°

$$\theta_0 = \theta_1, \theta_2 = 30^\circ$$

Why peak at z~1?



Why peak at z~1?

 $M \propto \delta r^3, F \propto -\delta r$ $\Psi \propto -\delta r^2 + const.$ $\frac{\Delta T}{T} \propto -\frac{\partial \Psi}{\partial \eta} \Delta \eta \propto -\delta r^3 \frac{d\Psi_t}{d\eta} d\eta$

Why peak at z~1?



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Significance of ISW contribution

$$P_{all} = P_{LSS} P_{ISW} b$$



ISW effect vs LSS effect



z	δ_0^L	$ heta_0$	$ heta_1$	r_0	$\Delta T_f(\text{void})$	$P_{\rm all}(\%)$
1.0	-0.0085	14°	30°	$570h^{-1}{ m Mpc}$	$7.7 \mu { m K}$	0.68

Why 12 % ISW contribution?

$$P = P_{LSS} P_{ISW} b \quad \Delta T = \Delta T_{LSS} + \Delta T_{ISW}$$

 $P_{\rm max} \sim P_{LSS} (\Delta T / 2) P_{ISW} (\Delta T / 2) b$

 $LSS: 20\,\mu\text{K} \times 3 = 60\,\mu\text{K}$ $void: 3\,\mu\text{K} \times 3 = 9\,\mu\text{K}$







Even in an accelerated epoch, a hot ring surrounding a cold spot can be produced if a void is locally over-compensated but it is less probable.

Assuming that a single moderately undercompensated low density region in the direction to the CS at z<1.5, the expected contribution to the CMB is about 10 % and the radius is about 600 h^{-1} Mpc and $\delta \sim -0.009$ at z~1. The probability of chance alignment is 0.7%. 100% LSS is 7 times less probable.







 Deviation from perfect spherical symmetry may be important as the assumed underdense region is in linear regime.

Other underdense regions in the l.o.s.may be also important especially if spatial correlation is important (void-in-void) (Tomita & Inoue `11 in preparation).







If a low density region at z~1 does not exist, we need to consider more anomalous fluctuations at z<1 or z>1.

Photo-z or spectroscopic survey aiming at z<1.5 with a wide-range field (>100 degree^2) to the CS is important.

About 10 % decrease in TE correlation in the l.o.s to the CS is expected (observable by Planck?)